# The Crystal Structure of Trimesic Acid (Benzene-1, 3,5-tricarboxylic Acid)* 

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The crystal structure of one modification of trimesic acid (benzene-1,3,5-tricarboxylic acid) has been determined. The crystals are monoclinic, space group $C 2 / c$, with $a=26 \cdot 52, b=16 \cdot 42, c=26 \cdot 55 \AA$, and $\beta=91.53^{\circ}$; there are 48 molecules in the cell and, hence, 6 in the asymmetric unit. Intensities were estimated visually from Weissenberg photographs; about 11,500 were recorded of which 3200 were too weak to be observed. The structure was derived by consideration of Patterson maps, optical transforms and molecular packing; refinement was by multiple-matrix least squares. In all, 979 parameters were refined, including anisotropic temperature facters for the heavy atoms and coordinates and isotropic temperature factors for the hydrogen atoms. The final $R$ index is 0.068 and the standard deviation of fit is 1.49 . The basic structural motif is a continuous, two-dimensional network comprising large, open rings formed by six molecules of trimesic acid bonded together through pairs of $\mathrm{O}-\mathrm{H} \cdots \mathrm{O}$ hydrogen bonds. The networks are not planar, pleats involving dihedral angles of about $70^{\circ}$ being formed by means of a twisting of scme of the hydrcgen-bonding carboxyl groups out of the planes of their benzene rings. The networks continually interpenetrate one another; through each large ring of one network pass three parallel rings of other networks. The average C(ring)-C(ring) distance is $1.390 \AA$, and the average $C$ (ring)-C(carboxyl) distance is $1.489 \AA$; nene of the individual values differs significantly from the average. On the other hand, there is a wide spectrum of $\mathrm{C}-\mathrm{O}$ distances. For the carboxyl groups that are severely twisted out of the plane of the benzene rings, the $\mathrm{C}-\mathrm{O}(\mathrm{H})$ and $\mathrm{C}=\mathrm{O}$ distances differ greatly, the average values being 1.319 and $1.218 \AA$; for the carboxyl groups that are only slightly twisted, the distances are mere nearly equal ( 1.279 and $1.245 \AA$ ). For a third category of carboxyl groups the hydrogen atoms are disordered, scme in ratios of $2: 1$ and others $1: 1$. In these, the $\mathrm{C}-\mathrm{O}$ distances reflect the disorder.

## Introduction

Our attention was first directed to trimesic acid - ben-zene-1,3,5-tricarboxylic acid, $\mathrm{C}_{6} \mathrm{H}_{3}(\mathrm{COOH})_{3}$ - by Professor H. McConnell and Dr G.R. Liebling. Their e.p.r. experiments indicated that trimesic acid may be damaged by X-radiation to produce a 'relatively high concentration of stable, long-lived free radicals' (Liebling, 1965); they also pointed out that knowledge of the crystal structure was essential to the precise interpretation of their observations.

## Experimental

Large prismatic crystals of trimesic acid were kindly furnished by Dr Liebling. Preliminary oscillation and Weissenberg photographs showed them to be monoclinic, space group Cc or C2/c ( $h k l$ absent for $h+k$ odd, $h 0 l$ absent for $l$ odd), with approximate cell dimensions $a=c=26 \cdot 5, b=16 \cdot 35 \AA, \beta=91 \cdot 6^{\circ}$. (More accurate values, obtained subsequently from a leastsquares treatment of measurements of $h 0 l$ and 0 kl Weissenberg photographs prepared in a Straumanistype camera, are given in Table 1.) The density meas-

[^0]ured by flotation is $1.46 \pm 0.02 \mathrm{~g} . \mathrm{cm}^{-3}$, indicating 48 molecules in the unit cell.*

## Table 1. Crystal data

Trimesic acid $\quad \mathrm{C}_{9} \mathrm{H}_{6} \mathrm{O}_{6}$
Formula weight $210 \cdot 14$
Monoclinic, $\quad$ Space group C2/c

$$
\begin{aligned}
& a=26.520 \pm 0.002 \AA \\
& b=16.420 \pm 0.001 \\
& c=26.551 \pm 0.002 \\
& \beta=91.53 \pm 0.01^{\circ}
\end{aligned}
$$

$$
\begin{array}{lr}
D_{x}=1 \cdot 449 \mathrm{~g} . \mathrm{cm}^{-3} & Z=48 \\
D_{m}=1.46 \mathrm{~g} \cdot \mathrm{~cm}^{-3} & F(000)=5184
\end{array}
$$

* Many of the crystals were found to be twinned about the (101) or (10T) planes. In an attempt to obtain additional untwinned crystals, some of the material was recrystallized from water. However, this resulted in crystals of a different modification - also monoclinic, space group $C c$ or $C 2 / c$, but with approximate cell dimensions $a=26.2, b=7.8, c=21.9 \AA$, $\beta=102^{\circ}$. On oscillation photographs about the $b$ (needle) axis, the odd layer lines were much weaker than the even; moreover, Weissenberg photographs of these odd levels showed the spots to be extensively smeared along the reciprocal lattice rows parallel to $\mathbf{c}^{*}$, indicating disorder. These crystals probably contain water, as the density calculated on the basis of 16 molecules of trimesic acid per unit cell is only $1.3 \mathrm{~g} . \mathrm{cm}^{-3}$. A study of the weighted reciprocal latice suggested that the rings are approximately perpendicular to the $b$ axis with one edge parallel to $c^{*}$. No further work on this modification was undertaken.

A number of preliminary observations encouraged us to proceed with the structure analysis. A rotation photograph about the $b$ axis (Fig. 1) showed a rather clear transform of a benzene ring, and suggested that at least some of the molecules are parallel to $\mathbf{b}$. The $h 0 l$ Weissenberg photographs had almost exact mm symmetry, the mirror planes being (101) and (101). (On upper-layer photographs this mm pseudosymmetry disappears.) Moreover, some extremely strong $h 0 l$ reflections indicated the orientations, in projection, of the rings; they also suggested the presence of a center of symmetry and, hence, that the space group is $C 2 / c$ rather than $C c$. For space group $C 2 / c$, the asymmetric unit contains six molecules of trimesic acid.

Complete intensity data were collected by the mul-tiple-film equi-inclination Weissenberg technique using $\mathrm{Cu} K \alpha$ radiation. For this purpose, two large, untwinned crystals were eventually found and ground into spheres by the method of Bond (1951). One, 0.43 mm in diameter, was mounted about the $b$ axis and layer lines $0-14$ were recorded; it was later remounted about a and layer lines $0-4$ were recorded. The second crystal, about the same size, accidently cleaved into two equal hemispheres during the grinding. It was mounted perpendicular to the cleavage plane (101) - that is, parallel to the [ 10 T ] axis - and the zero level was recorded in order to permit correlation of the data from the other axes (because of the $C$ centering, the odd layers about b have no reflections in common with the even layers about a, and vice versa). The intensities were estimated visually. Approximately 11,000 reflections were recorded on the photographs about the $b$ axis, 2500 about a, and 200 about [101]; in all, 11,563 independent reflections were measured, of which 3172 were too weak to be observed. The reciprocal sphere for copper radiation contains approximately 12,500 reflections.

The data reduction, as well as all subsequent calculations, was carried out on an IBM 7094 computer using subroutines of the CRYRM system (Duchamp, 1964). The intensity values as estimated on each film (and which ranged from about 2 to 65 for the comparison strip we used) were read into the computer and each was immediately assigned a standard deviation according to the function

$$
\begin{align*}
\sigma(I)=(0 \cdot 833+m I & \left.\frac{0 \cdot 1 I^{2}}{(65-I)^{2}}\right)\{1+0 \cdot 25 \\
& \left.\times \exp \left[-50\left(0 \cdot 5-\sin ^{2} \theta\right)^{2}\right]\right\} \tag{1}
\end{align*}
$$

Within the first factor, the constant term increases the uncertainties when $I$ is small and the quadratic term increases the uncertainties when $I$ is large (if $I$ is equal to or greater than the limiting value $65, \sigma$ is set to a very large value and the weight of that observation becomes essentially zero); the second factor increases the uncertainties of reflections occurring in the region
where the $\alpha_{1}-\alpha_{2}$ splitting begins. The linear coefficient $m$ was at first assigned the value $0 \cdot 1$.

These standard deviations were used in the calculation of the intensity factors relating different films of a particular layer. Empirical values of $m$ (equation 1) were also calculated at this time, from a comparison of different measurements of the same reflections. The values were 0.043 for data collected about the $a$ axis, 0.059 for data about $\mathbf{b}$, and 0.166 for data about [101] ]. New standard deviations were assigned, based on these empirical values of $m$; these standard deviations were propagated through all subsequent corrections and scaling for use later in least-squares weighting. All the data within each layer were then scaled together, averaged, and multiplied by 1/Lp. No absorption correction seemed necessary ( $\mu t \simeq 0 \cdot 5$ ).

The final step of the scaling procedure was the determination of the scale factors ( $C_{i} \cdots n$ ) relating the various sets of films. These factors were calculated by a method which minimizes $\Sigma w_{i}\left(C_{k} I_{i k}-C_{l} I_{i}\right)^{2}$, the sum-


Fig.1. A rotation photograph $(\mathrm{Cu} K \alpha)$ of trimesic acid. The rotation axis (b) is horizontal.
mation being over all reflections $i$ common to film sets $k$ and $l$. Using these factors, an averaged intensity and standard deviation was assigned to each of the 11,563 independent reflections.

## Determination of the structure

Several different methods of attack were used in deriving a trial structure. Perhaps the biggest clue came from comparing the Fourier transform of the molecule (synthesized optically by means of an apparatus constructed by Professor J.H.Sturdivant and making use of two $12^{\prime \prime}$ spherical lenses) with zero-level precession photographs prepared at $15^{\circ}$ intervals about the $b$ axis. This comparison confirmed the previous observation that at least some molecules are parallel to $\mathbf{b}$, and suggested three possible orientations: approximately perpendicular to a, to $\mathbf{c}$, and to ( $10 \overline{\mathrm{I}}$ ).

In view of the length of the $b$ axis and the observed orientations of some of the rings, it now appeared probable that one of the features of the structure was the formation of planar ribbons of molecules, hydrogen bonded together through their COOH groups; these ribbons presumably were parallel to $b$ and in the orientation suggested by the molecular transform work. When viewed down the $b$ axis, many of the atoms within these ribbons would be closely superimposed, and hence Patterson projections on to (010) were expected to have many of the advantages of a heavyatom map. Sharpened and unsharpened Patterson projections were calculated; and while they differed from one another in many important features, both were in agreement with the proposed ribbon-like arrangement.

Consideration of possible intermolecular arrangements as viewed down the $b$ axis eventually led to a satisfactory structure in projection. The planar ribbon hypothesis and the possible orientations deduced from the molecular transform work were used as a starting point. The Patterson maps were used to check each step and to give hints on how to proceed. The threedimensional implications for each possible arrangement were also studied. The only arrangement satisfying all criteria was one in which all six molecules in the asymmetric unit were parallel to the $b$ axis, three being approximately perpendicular to $\mathbf{a}$ and three to $\mathbf{c}$.

Three cycles of full-matrix least-squares refinement of the $x$ and $z$ coordinates were calculated, using only the $h 0 l$ data. Even though this refinement proved to be rather unstable, necessitating careful evaluation of the indicated shifts, the $R$ index for the $h 0 l$ reflections was reduced from 0.61 to 0.38 by this method. Significant shifts in the centers of the molecules resulted, as well as twists of 4 of the 18 carboxyl groups out of the planes of their benzene rings.

The $y$ coordinates of the centers of the six molecules were derived from a sharpened three-dimensional Patterson map. Coordinates for all the 90 heavier atoms in the asymmetric unit were then calculated assuming standard bond distances and angles, and structure fac-
tors were calculated for the complete three-dimensional data set; the $R$ index ( $R=\Sigma\left\|F_{o}|-s| F_{c}\right\| / \Sigma\left|F_{o}\right|$ ) was $0 \cdot 51$.

## Three-dimensional refinement

Refinement was primarily by least squares, the quantity minimized being $\Sigma w\left(F_{o}^{2}-s^{2} F_{c}^{2}\right)^{2}$ (where $1 / s$ is the scale factor). Atomic scattering factors were taken from International Tables for X-ray Crystallography (1962). There was one false start: the structure derived from the Patterson maps required that some of the carboxyl groups of structurally equivalent molecules be hydrogen bonded to one another across twofold axes, and hence that the hydrogen atoms be disordered. An alternate ordered structure with hydrogen bonding across centers of symmetry (and identical in projection down b) could be generated by translating the origin to $x=\frac{1}{4}$, $z=\frac{1}{4}$. Initial refinement was based on this alternate structure. However, reflections with $(k+l)$ odd showed continuing poor agreement between observed and calculated structure factors; the agreement improved markedly when the original structure was considered.

The amount of core storage available limited the number of parameters that could be included in a single matrix to approximately 200 , and accordingly the parameters being refined were distributed among a number of smaller matrices. In the early stages, when only the coordinates and isotropic temperature factors of the carbon and oxygen atoms were being adjusted, the coordinates of the atoms in four molecules ( 180 parameters) could be combined in one matrix, the remaining parameters being held fixed; alternatively, four matrices could be collected: three, of 90 parameters, each containing the coordinates of two molecules and the fourth, of 91 parameters, containing the scale and temperature factors. Other set-ups were also used.

Initial calculations were based on the 1070 reflections with $\sin ^{2} 0 / \hat{\lambda}^{2}<0 \cdot 08$, and with the waighting scheme $w=1 / F_{o}^{2}$ for $F_{o} \geq 15, w=1 / 15 F_{o}$ for $F_{o} \leq 15$. After several cycles of refining only the coordinates, refinement was speeded up by applying the least-squares method of Waser (1963), in which constraints - in the form of additional, heavily weighted observational equations were introduced to insure that the molecules retained their expected geometry. In these calculations the benzene rings and the carboxyl groups were required to be planar, but not necessarily coplanar. The coordinates and isotropic temperature factors were then refined using an increased amount of data and weights equal to $1 / \sigma^{2}\left(F_{o}^{2}\right)$. Eventually the $R$ index for the complete set of data was reduced to $0 \cdot 16$.

At this stage anisotropic thermal parameters for the 90 heavy atoms were introduced. For the carbon atoms, the starting values were the isotropic parameters from the preceding cycles; since the thermal motions of the oxygen atoms were expected to be quite anisotropic (larger than average in directions perpendicular to the planes of the carboxyl groups and smaller than average along the $\mathrm{C}-\mathrm{O}$ bonds), their starting values were esti-
mated anisotropic coefficients. The complete set of data ( 11,563 reflections) was included in the remaining refinement cycles, and weights were taken equal to 1/ $\sigma^{2}\left(F_{o}^{2}\right)$ (see Experimental). Since the number of parameters was now up to 811 ( 270 coordinates, 540 temperature parameters, and a scale factor), further subdivision of the matrices was necessary. We also found it efficient to carry out the least-squares refinement cycles in pairs: in the second cycle only the shift vectors were calculated, the inverse matrices of the normal equations being saved from the first cycle.

Four cycles of block-matrix refinement (two matrices, one for the coordinates and the other for the $b_{i j}$ 's, for each atom) reduced the $R$ index to $0 \cdot 10$. However, the molecular geometry was not entirely satisfactory, and it was apparent that convergence was being prevented by a small number - less than a dozen of reflections with large residuals $V^{\prime} w\left(F_{o}^{2}-s F_{c}^{2}\right)$, and that in most cases the residuals were probably due to the neglect of the hydrogen atoms. A three-dimensional difference map was therefore calculated; it indicated unambiguously the positions of the ring hydrogen
atoms, but not of the carboxyl hydrogen atoms. This map had a rather high noise level, undoubtedly due to incorrect parameters of the heavier atoms.

Nine more refinement cycles, with the ring hydrogen atoms included only in the structure factors, led to an $R$ index of 0.079 . Once again, several matrix arrangements were tried; one of the most efficient procedures seemed to be to include the coordinates of all the atoms of one molecule in a single matrix at the expense of reducing the sizes of the temperature-factor matrices. However, satisfactory convergence was still not attained.

Further attention was now given to the problem of locating the hydrogen atoms of the carboxyl groups. Difference maps were calculated through the ten sets of four oxygen atoms involved in the hydrogen bonding between carboxyl groups. These maps clearly showed the positions of two of the three hydrogen atoms [ $\mathrm{H}(5)$ and $\mathrm{H}(6)$ ] of each molecule; the third hydrogen atom $[\mathrm{H}(4)]$ of each molecule appeared to be disordered. For molecules $A$ and $B$, the hydrogen bonding is across a twofold axis and accordingly $\mathrm{H}(4)$


Fig. 2. Composite electron-density maps calculated at the conclusion of the refinement. Top: Sections evaluated in the planes of the benzene rings and, where necessary, in parallel planes containing carboxyl oxygen atoms. Contours are at intervals of $2 \mathrm{e} . \AA^{-3}$ beginning with $1.0 \mathrm{e} . \AA^{-3}$ (dashed). Bottom: Similar sections of difference maps, the contributions of the hydrogen atoms being omitted from the $F_{c}$ 's. Contours are $0 \cdot 1$ (dashed), $0 \cdot 2,0 \cdot 3 \ldots$ e. $\AA^{-3}$. The $b$ axis is vertical; for the left half of the Figure the view is roughly along $c$ while for the right half it is roughly alonga. Thus there is a fold (not shown) of about $70^{\circ}$ between the plane of molecules $A, C$, and $E$ and that of $B, D$, and $F$.
must be disordered in the form of two half-atoms. For molecules $C, D, E$, and $F, \mathbf{H}(4)$ appeared to be only partially disordered, one of the two positions being favored over the other. The apparent disorder of these four atoms was tentatively ignored, each atom being placed in the more favorable of the two positions.
Eight more least-squares cycles were now calculated. In the first four, the coordinates of the hydrogen atoms were included in a single ( $114 \times 114$ ) matrix; in the next four, isotropic temperature factors of the hydrogen atoms were included in an additional $38 \times 38$ matrix. During these cycles, the hydrogen atoms $\mathrm{H}(4)$ for molecules $C, D, E$, and $F$ shifted to positions approximately midway between the two hydrogen-bonding oxygen atoms. Another set of difference maps was therefore calculated for these atoms. These maps strongly suggested that the four hydrogen atoms are disordered in a ratio of about 2 to 1 ; that is, it appeared as though $\frac{2}{3}$ of an atom should be placed in one site (in each case, the position 'favored' in the preceding map) and $\frac{1}{3}$ in the alternate site adjacent to the other oxygen atom.

Two more least-squares cycles completed the refinement. A total of 979 parameters were adjusted: three coordinates and six temperature factors for each of 90 heavy atoms, three coordinates and one temperature factor for each of 42 hydrogen atoms (five whole atoms and one split atom for each of the six molecules), and a scale factor. Calculation time on an IBM 7094 computer was approximately 2 hours for the first cycle and 15 minutes for the second (in which only the structure factors and the shift vectors were calculated).

During the last cycle, the shifts in the heavy-atom parameters were small - almost invariably less than $0 \cdot 1 \sigma$. Three coordinates and one temperature factor of disordered hydrogen atoms shifted by approximately $1 \cdot 5 \sigma$. The final $R$ index is 0.068 ; the 'goodness of fit' $\left(=\left[\Sigma n\left(F_{o}^{2}-s^{2} F_{c}^{2}\right)^{2} /(m-s)\right]^{1 / 2}\right)$ is $1 \cdot 49$. Of 3172 reflections too weak to be observed, 86 calculated larger than the threshold value.

In Fig. 2 are shown electron-density and difference maps calculated at the conclusion of the refinement. In the upper half is shown a composite of a number of electron-density sections; these were calculated in the least-squares planes of the carbon atoms of each molecule except for those oxygen atoms which lie more than $0 \cdot 20 \AA$ from the plane, in which case a parallel plane through the center of the oxygen atom was calculated. In the bottom half is shown a composite of difference sections calculated with the hydrogen atoms omitted from the structure factors; for the ring hydrogen atoms these sections were calculated in the planes of the benzene rings, while for the carboxyl hydrogen atoms the sections are in the least-squares planes of the four oxygen atoms involved in hydrogen bonding between pairs of carboxyl groups.

Observed and calculated structure factors are listed in Table 2; the final atomic parameters and their standard deviations are listed in Table 3. The standard deviations were calculated in the usual way from the
residuals and the diagonal elements of the inverse matrices of the normal equations; since these matrices were incomplete, the standard deviations are presumably underestimated. However, the agreement between the many different measurements of chemically equivalent bond distances suggests that this underestimation is not serious (see Discussion).

## Discussion of the results

## Bond distances and angles

The bond distances and most of the bond angles are shown in Fig. 3, which is a representation of a portion of the structure viewed normal to the planes of the molecules except that a fold, running vertically between molecules of types $E$ and $F$ and involving a dihedral angle of approximately $70^{\circ}$, is not shown. The interior angles of the benzene rings are not given in this Figure; they range from $119 \cdot 2^{\circ}$ to $121 \cdot 0^{\circ}$, and there seems to be no systematic difference between angles opposite carboxyl groups and those opposite hydrogen atoms.

The standard deviations in positional parameters (Table 3) correspond to e.s.d.'s in the $\mathrm{C}-\mathrm{C}$ and $\mathrm{C}-\mathrm{O}$ distances of about $0.004 \AA$ and of about $0.4^{\circ}$ in the angles. There are 36 structurally distinct ring $\mathrm{C}-\mathrm{C}$ distances, all of which may be expected on chemical grounds to be equal. The scatter of these 36 measurements from their average value of $1 \cdot 390 \AA$ leads to an internal standard deviation of $0.0058 \AA( \pm 0.0007 \AA)$; similarly, the scatter among the internal bond angles leads to a standard deviation of $0.43^{\circ}\left( \pm 0.05^{\circ}\right)$. We feel that this agreement between the two sets of standard deviations is satisfactory, and suggests that the refinement has been carried essentially to completion.

Least-squares planes were calculated through the six ring atoms of each molecule (with all six atoms weighted equally). Equations for these planes are given in Table 4, together with the out-of-plane distances of the various atoms in each molecule. Each group of six ring atoms of molecules $A, B, C$, and $D$ is probably coplanar within experimental error; for molecules $E$ and $F$ they are slightly, but significantly, not coplanar. None of the carboxyl groups is coplanar with its benzene ring; the largest deviations are for the carboxyl groups $\mathrm{C}(8)$ and $\mathrm{C}(9)$ of molecules $E$ and $F$, which are twisted about the C(ring)-C(carboxyl) bonds by amounts ranging up to $27^{\circ}$. These twists correspond to the folds in the hydrogen-bonded network, as will be discussed later. Besides being twisted about the C-C bonds, the carboxyl groups are also displaced through a bending of these bonds out of the planes of the benzene rings.

The $\mathrm{C}-\mathrm{O}$ distances suggest three distinct types of carboxyl group: those involving $\mathrm{C}(8)$ and $\mathrm{C}(9)$ of molecules $A, B, C$, and $D$; those involving $\mathrm{C}(8)$ and $\mathrm{C}(9)$ of molecules $E$ and $F$; and those involving $C(7)$ of all six molecules. The carboxyl groups of the first type are approximately coplanar with their benzene rings; for them, the two C-Odistances are not greatly different (av-

Table 2．Observed and calculated structure factors
The three columns contain values，reading from left to right，of $h, F$（obs）and $F(c a l)$ ．A minus sign preceding $F$（obs）should be read＇less than＇；asterisks indicate reflections omitted from the refinement．




























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Table 2 (cont.)




































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Table 2 (cont.)
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erage $\mathrm{C}=\mathrm{O}, 1 \cdot 245 \AA ; \mathrm{C}-\mathrm{O}(\mathrm{H}), 1 \cdot 279 \AA$ ), suggesting considerable resonance. The carboxyl groups of the second type are twisted out of the planes of the benzenerings; for them, the $\mathrm{C}-\mathrm{O}$ distances ( 1.218 and $1.319 \AA$ ) suggest that the double bond is much more localized. Finally, the carboxyl groups of $\mathrm{C}(7)$ all involve disordered hydrogen atoms; the $\mathrm{C}-\mathrm{O}$ distances reflect the disorder, being essentially equal in length for molecules $A$ and $B$, where the two disordered structures contribute equally, but significantly different (and in the appropriate sense) for the remaining molecules where one of the two structures is approximately twice as important as the other.

There is no apparent variation in the C(ring)-C(carboxyl) distance, none of the individual values differing by more than $0.014 \AA$ from the average value of $1.489 \AA$. The $\mathrm{O} \cdots \mathrm{O}$ hydrogen bond distances have values ranging from $2 \cdot 605 \AA$ to $2 \cdot 699 \AA$; the variations do not appear to be correlated with the type of carboxyl group (see above).

The apparent $\mathrm{C}-\mathrm{H}$ distances range from 0.79 to $1.03 \AA$, and the $\mathrm{O}-\mathrm{H}$ distances from 0.82 to $1.27 \AA$
(the latter value involving a disordered hydrogen atom and hence being of low precision).

## The temperature factors

The magnitudes and direction cosines of the principal axes of the temperature-factor ellipsoids are given in Table 5. The most pronounced anisotropies are associated with the oxygen atoms and, as anticipated, involve large out-of-plane displacements. The magnitudes of the displacements suggest that the values reported for some of the $\mathrm{C}-\mathrm{O}$ bond distances should be increased by amounts ranging up to $0.03 \AA$. We have not made such corrections, as the resulting $\mathrm{C}-\mathrm{O}$ distances appear to be less reasonable than the uncorrected values.

## The hydrogen bonding

The most interesting features of the structure involve the intricate arrangement of hydrogen bonds and the resulting three-dimensional network. A portion of this network is shown in the stereographic photographs Fig.4(b)-(f).
Table 3. The final parameters and their standard deviations (in parentheses)

The coordinates of the heavy atoms have been multiplied by $10^{4}$; those of the hydrogen atoms, by 103 . The anisotropic temperature parameters of the heavy atoms are of the form $\exp \left[-\left(b_{11} h^{2}+b_{22} k^{2}+b_{33} l^{2}+b_{12} h k+b_{13} h l+b_{23} k l\right)\right]$, and have been multiplied by $10^{5}$. For the hydrogen atoms the values listed are of the isotropic temperature










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Table 4. Least-squares planes through atoms $\mathrm{C}(1), \ldots, \mathrm{C}(6)$



Fig. 3. The bond distances and angles for one asymmetric unit. The interior angles of the benzene rings are not given; they range from $119 \cdot 2^{\circ}$ to $121 \cdot 0^{\circ}$. The view is the same as for Fig. 2

(a)

(b)

(c)

Fig. 4. Photographs of a model of the structure. (a) The trimesic acid molecule. (b) . . (f) Stereographic photographs of portions of the three-dimensional structure. (b) A portion of a single 'chicken-wire' network, showing the large six-molecule rings and the fold in the network. (c) Two interlocking networks.

(d)

(e)

(f)

Fig. 4 (cont.) (d) Three networks. (e) Four networks. Here can be seen the interpenetration of one network by parallel portions of three other networks. $(f)$ Six networks, showing the mutual interlocking of three parallel networks with three others in a different orientation.

The basic motif is that shown in Fig. 5 - an open framework, much like chicken wire, comprising rings of six molecules of trimesic acid hydrogen bonded together through pairs of $\mathrm{O}-\mathrm{H} \cdots \mathrm{O}$ hydrogen bonds between carboxyl groups. A brief consideration of molecular models makes it clear that this motif is the only reasonable array that can be formed by planar trimesic acid molecules; however, it results in large holes in the centers of the six-molecule rings, and these holes must be filled if satisfactory packing is to be attained. The filling is accomplished through a mutual interpenetration of different chicken-wire frameworks. Through each hole of one framework pass, at an angle of about $70^{\circ}$, three parallel six-molecule rings of other frameworks; in turn, the original framework is one of three arrays that pass through the holes in the second set of frameworks.

The details of this interpenetration are shown in Fig.6, which is a schematic view down the $b$ axis. Mol-
ecules of type $A, C$, and $E$ are hydrogen-bonded to one another, as shown in Figs. 2 and 3, to form three parallel, planar arrays approximately perpendicular to the $c$ axis, while molecules $B, D$, and $F$ form similar arrays approximately perpendicular to a. These arrays interpenetrate one another. Moreover, hydrogen bonds between molecules of types $E$ and $F$ connect one array directly to the other, this connection involving a dihedral angle of approximately $70^{\circ}$; thus, each chicken-wire array is folded into pleats. The pleating occurs between carboxyl groups $\mathrm{C}(8)$ and $\mathrm{C}(9)$ of molecules of types $E$ and $F$, and is accomplished by means of the large out-of-plane twists of these carboxyl groups from the planes of their benzene rings.

The dihedral angles between planes of adjacent molecules within a three-molecule stack are $2.6^{\circ}$ for $A$ and $E, 3.8^{\circ}$ for $A$ and $C^{\prime}$ (related to $C$ by a $2_{1}$ axis), $1.6^{\circ}$ for $F$ and $B$, and $2.5^{\circ}$ for $B$ and $D^{\prime}$

Table 5. Principal axes of the thermal ellipsoids
Columns are the magnitudes, $B$, and the direction cosines ( $\times 10^{3}$ ) wihh respect to the crystallographic axes, $a, b$, and $c$ respectively.

|  | A | B | C | D |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  |  |  |  |  |
| c(2) |  <br>  |  |  |  |  |  |
| c(3) |  |  |  |  |  |  |
| c(4) |  | (ex |  |  $2.42{ }_{265} 26$ |  |  |
| c(s) |  |  |  |  |  |  |
| c(6) |  |  |  |  |  | (2, |
|  | $\begin{array}{lll} 3.77 \\ \hline \end{array}$ |  |  |  |  | (ellll |
| ${ }^{\text {c } 8)}$ |  |  | ¢ |  | $\begin{array}{lll} 3.10 & -817 \\ \hline \end{array}$ |  |
| c(9) |  | (1) |  |  |  |  |
|  | 10.16 -106 des 98 <br>  |  |  |  | (1) |  |
| (2) |  $\begin{array}{llll}1.22 & 281 & 855 & 800 \\ 2.33 & 912 & -97 & 87\end{array}$ |  |  |  |  |  |
| (3) |  |  $\begin{array}{llll}3.26 & 400 & -118 & 688 \\ 2.67 & 099 & 993 & 111\end{array}$ |  |  |  |  |
| $\alpha 4$ |  |  |  |  | (ex |  |
| 0(3) |  |  |  |  |  |  |
| (6) |  |  |  |  |  |  |

(related to $D$ by a center of symmetry). The average distance between molecules within a stack is about $3 \cdot 1 \AA$.

Pairs of molecules of type $A$ are hydrogen bonded to one another across a twofold axis (see Fig.6), as are molecules of type $B$. The hydrogen atoms involved in these bonds must, then, be disordered (the $\mathrm{O} \cdots \mathrm{O}$ distances are far too long to represent symmetric hydrogen bonds). It is interesting that the other hydrogen bonds oriented perpendicular to the $b$ axis - those between carboxyl groups $C(7)$ of molecules $C$ and $E$ and of molecules $D$ and $F$ - also appear to be disordered, although in these cases one of the two arrangements seems to be favored over the other. The remaining hydrogen bonds - those between carboxyl groups $C(8)$ and $C(9)$ of different molecules - show no such disorder. Moreover, each of the twelve $C(8)-C(9)$ carboxyl groups is oriented such that its hydrogen atom is at the maximum distance along $b$ from the center of its respective molecule. The result of this systematic arrangement of hydrogen bonds is that the holes in the centers of the six-molecule rings are slightly ellip-


Fig. 5. The basic chicken-wire motif - a two-dimensional network of six-molecule rings. Dashed lines represent hydrogen bonds.


Fig.6. A schematic representation of the structure viewed down the $b$ axis. Individual trimesic acid molecules are represented by straight, solid lines. Dotted lines represent hydrogen bonds between $n$ olecules having sin ilar $y$ coordinates; these horizontal hydrogen bonds all involve disordered hydrogen atoms (see Figs. 2 and 3). The pairs of curved lines represent hydrogen bonds between molecules lying at different heights along $\mathbf{b}$; for these, the hydrogen atoms seem to be entirely ordered. The molecules labelled $A$ through $F$ are those whose coordinates are listed in Table 3.
tical, the diameter being greater perpendicular to the $b$ axis than parallel to $\mathbf{b}$.

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# The Crystal Structures of Copper Tetrammine Complexes A. $\mathrm{Cu}\left(\mathrm{NH}_{3}\right)_{4}^{4} \mathrm{SO}_{4} \cdot \mathrm{H}_{2} \mathrm{O}$ and $\mathrm{Cu}\left(\mathrm{NH}_{3}\right)_{4} \mathrm{SeO}_{4}{ }^{*}$ 

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#### Abstract

The room temperature crystal structures of two copper tetrammine complexes have been refined by the full-matrix least-squares method using three-dimensional Mo $K x$ intensity data. The space group for $\mathrm{Cu}\left(\mathrm{NH}_{3}\right)_{4} \mathrm{SO}_{4} . \mathrm{H}_{2} \mathrm{O}$ is Pnam with lattice constants $a_{0}=10.651, b_{0}=11.986$, and $c_{0}=7.069 \AA$; that for $\mathrm{Cu}\left(\mathrm{NH}_{3}\right)_{4} \mathrm{SeO}_{4}$ is $P 2_{1} / n$ with lattice constants $a_{0}=10 \cdot 313, b_{0}=10 \cdot 259, c_{0}=7 \cdot 405 \AA$ and $\beta=104.43^{\circ}$. In both complexes, the copper ion is surrounded by four ammonia molecules ( $\sim 2.0 \AA$ ) in a near square planar arrangement and by two more distant oxygen atoms (but at unequal distances from the copper ion). The copper ions are connected via the oxygen atoms in a chain-like arrangement, and a network of hydrogen bonds hold these chains and ions together.


## Introduction

The magnetic and thermal properties of $\mathrm{Cu}\left(\mathrm{NH}_{3}\right)_{4} \mathrm{SO}_{4} \cdot \mathrm{H}_{2} \mathrm{O}$ (CTASUL) have been extensively studied and have been interpreted as those of a substance containing a magnetic linear chain structure(Saito \& Kanda, 1967; Rogers, Carboni \& Richards, 1967). Recently Saito \& Kanda (1967) conducted proton nuclear magnetic resonance studies on this material in the paramagnetic and ordered state. Unfortunately the lack of proton positions prevented a quantitative interpretation of these data. Since the early crystallographic study by Mazzi (1955)) reported a structure based on projected data (with overlapping atomic sites) and indicated an unusual coordination about the copper ion, a detailed structure determination was considered necessary for any future quantitative interpretation of the observed resonance data.

Our attempts to crystallize an isostructural selenate compound proved to be unsuccessful; however, we isolated several other copper tetrammine complexes. The crystal structure of one of these complexes, $\mathrm{Cu}\left(\mathrm{NH}_{3}\right)_{4} \mathrm{SeO}_{4}$ (CTASEL), was determined and the coordination about the copper was shown to be distorted in a similar manner to that found in CTASUL. The room temperature X-ray diffraction studies on these two compounds reported here are part of a program to gain an understanding of the relationship be-

[^1]tween chemical bonding and magnetic properties in materials containing linear chain structures.

## Experimental

The deep blue crystals of both compounds were grown in the interface between an alcohol layer and an aqueous solution, prepared by adding an excess of $\mathrm{NH}_{3}$ to the required copper salt solution. Crystals were examined by X-ray photographic techniques to insure that a single-crystal specimen rather than a twin was selected for our measurements. The space groups were selected on the basis of the observed extinctions on Mo $K \alpha$ precession photographs and the absence of a piezoelectric response. In addition, the statistics of the normal structure factors calculated from the intensity data on CTASEL corroborated a centrosymmetric space group. The final agreement between the observed and calculated structure factors indicate that the space groups selected are the correct choice. The cell edges have been selectcd to conform to Donnay's (1943) rules, and hence, may not agree with other published results. Specific values found from our experimental procedures are summarized in the following section.

The lattice constants for these complexes were obtained with $\mathrm{Cu} K \alpha$ radiation ( $\lambda$ for $K \alpha_{1}=1.54050 \AA$ ) by least-squares fit of high 20 values measured on films taken with a 115 mm diameter Weissenberg camera utilizing Straumanis film loading.

Three-dimensional Mo $K x$ intensity data were measured using the $\theta-20$ scan technique on a Picker dif-


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